PREPARAREA ELECTROLITULUI SOLID DE TIP (La,Sr)(Ga,Mg)O₃
PRINTR-O METODĂ ÎMБUNĂTĂȚITĂ A REACȚIILOR ÎN FAZĂ SOLIDĂ
PROCESSING OF (La,Sr)(Ga,Mg)O₃ SOLID ELECTROLYTE USING
AN ENHANCED SOLID STATE TECHNIQUE

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Sw and Mg doped lanthanum gallate perovskites (Laₓ,ₜ-Sr,ₜ-Gaₓ-Mgᵧ,ₜ-Oₙ-δ) shortened as LSGM are promising electrolyte materials for intermediate temperature solid oxide fuel cells (IT-SOFCs).

The evolution of secondary phases, such as LaSrGaO₃ has been identified to be a problem in the preparation of LSGM. Actually, no matter what technical preparation route was adopted (wet-based or solid state-based method) the presence of the secondary phases, accompanying the main phase LSGM was reported. The investigated by means of X-ray diffractometry, FT-IR analyses (DTA/TG) on the stoichiometric raw mixture and dilatometric measurements on sintered bodies were characterized using different techniques. Thermal treatment was reported in this work. The obtained ceramics pure Sr and Mg doped LaGaO₃ can be obtained.

Choosing an appropriate annealing program, almost pure Sr and Mg doped LaGaO₃ can be obtained.

Perovskite of tip lantanat de galiu dopați cu Sr și Mg (Laₓ,ₜ-Srₓ,ₜ-Gaₓ-Mgᵧ,ₜ-Oₙ-δ), abreviat LSGM, sunt electroli promițători pentru pîelele de combustie cu electroli solid de temperaturi intermediare (IT-SOFC).

1. Introduction

Perovskite-type ABO₃ phases derived from lanthanum gallate, LaGaO₃, possess a higher ionic conductivity than that of stabilized zirconia. Sr and Mg doped lanthanum gallate perovskites (Laₓ,ₜ-Srₓ,ₜ-Gaₓ-Mgᵧ,ₜ-Oₙ-δ), shortened as LSGM-X, where X and Y are the doping levels in mole percentage (mol%) at the La- or A-site and the Ga- or B-site, respectively, are promising electrolyte materials for intermediate temperature solid oxide fuel cells (IT-SOFCs) [1-11]. Compared with CeO₂-based electrolytes (which are also of interest for lowering the operating temperature of SOFCs), the electrolytic domain of doped LaGaO₃ extends to substantially lower oxygen chemical potentials [12-14]. For the LSGM series, the maximum ionic transport is achieved at x=0.10–0.20 and y=0.15–0.20 [15]; further acceptor-type doping leads to progressive vacancy association processes.

However, the concentration of transition metal dopants should be limited to below 3–7% as further additions lead to increasing electronic and decreasing ionic conductivities [16].

Disadvantages of LaGaO₃-based materials include possible reduction and volatilization of gallium oxide, formation of stable secondary phases during the processing, the relatively high cost of gallium and significant reactivity with perovskite electrodes under oxidizing conditions as well as with metal anodes in reducing conditions. These problems may be, to some extent, suppressed by the optimization of processing techniques.

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LSMG were signaled in literature [2,3,5,10,11]. In this work one obtain, by solid-state technique, a pure Sr and Mg doped LaGaO$_3$ ceramics using a specific annealing treatment.

2. Experimental and Characterization

Solid-state reaction (SSR) was chosen as synthesis route as it is the most commonly used technique in ceramic preparations. High pure compounds La$_2$O$_3$, SrCO$_3$, Ga$_2$O$_3$, 4MgCO$_3$.Mg(OH)$_2$.5H$_2$O were used as starting components. Stoichiometric amounts (La$_{0.8}$Sr$_{0.2}$Ga$_0.8$Mg$_0.2$O$_{3-x}$) were mixed in agate mortar, followed by the calcinations and sintering.

The green powder was first calcined at 1000°C for 6 h. A heating rate of 5°C/min was adopted up to 600°C and then 1°C/min up to 1000°C was used. The calcined powders were pressed uni-axially into cylindrical shape samples (ϕ=10 mm, h=2-3 mm). The compacts were then sintered in air box furnace at 1450°C and 1500°C respectively, for 6h.

The densities of sintered samples at 1500°C/6h (ρ~5.96 g/cm$^3$) and porosity (~1.0 %) were measured using the standard Archimedes method.

The structure of the obtained materials was determined by X-ray diffraction (Rigaku ULTIMA IV Diffractometer, CuK$_\alpha$). The search for impurity phases included the following substances: LaSrGaO$_4$ (Powder Diffraction File (PDF) No. 24–1208), SrLaGa$_2$O$_6$ (PDF No. 45-0637), La$_4$Ga$_2$O$_9$ (PDF No. 37–1433), SrGa$_2$O$_4$ (PDF No. 72-0222), La$_2$O$_2$.($CO_3$)$_{0.85}$ (PDF No. 84-1964).

IR (infrared) spectra were recorded using a Thermo NICOLET FTIR 6700 spectrometer. Thermal analysis (DTA/TG) was performed with a METTLER Toledo 851$^a$ Derivatograph. Thermal expansion coefficients (TEC) were measured using a Netzsch DIL 402 C dilatometer, constant heating rate of 5°C/min, in N$_2$. The morphology was determined by Scanning Electron Microscopy SEM (Zeiss DSM 942). The electrical behavior of the ceramic was characterized by impedance spectroscopy (Wayne Kerr 6440A) at frequencies in the range 20 Hz-3MHz with V$_{ac}$=1V. The measurements were carried out in air from room temperature up to 500°C. The pellets were coated with silver paste on both surfaces and then were heated at 300°C for 1 hour for a good contact with the electrode.

3. Results and discussion.

It is difficult to remove impurity phases (such as LaSrGa$_2$O$_4$) in the synthesis process of the Sr- and Mg-doped LaGaO$_3$, no matter what technique is employed. In fact, synthesis of a pure single phase of LSGM is rather hard to achieve [17]. Previous works concluded that, independently of preparation routes, calcination up to 1500°C is necessary to obtain a pure perovskite phase [18], because LaSrGa$_2$O$_7$ and less LaSrGaO$_4$ are in thermodynamic equilibrium with LSGM at 1400°C in air [19,20]. Another reason is the narrow composition range for the stability of the perovskite phase [21]. The solubility of Sr$^{2+}$ on the La$^{3+}$ sites in LaGaO$_3$ system is below 5 % for a sintering temperature of 1500°C. The solubility limit of Sr$^{2+}$ on the La$^{3+}$ sites increases by substituting Mg$^{2+}$ on the Ga$^{3+}$ sites in LaGaO$_3$ [22]. The results of Ishihara et al. [9] showed also that incorporation of the guest ions in the host lattice at specific sites is realized by substitution mechanism. The bulk diffusion coefficients are similar for all cations with activation energies which are strongly dependent on temperature. It was found that diffusion takes place through the bulk and along the grain boundaries [23,24].

Taking into consideration the above observations one can conclude that reactions taking place in the subsolidus domain in the corresponding phase diagram of the doped La$_2$O$_3$-Ga$_2$O$_3$ system may play a significant role in preventing (hindering) the parasite phase formation. Starting with carbonates as raw materials may constitute an advantage. The alkaline earth carbonates, MCO$_3$ (M = Ca, Sr, and Ba), have several interesting properties. At certain temperatures, the MCO$_3$ would transform polymorphically the structure by the rotational disorder transition of the anion group at a certain temperature, such as the transformation of CaCO$_3$ from orthorhombic to hexagonal at 470°C, SrCO$_3$ from orthorhombic to hexagonal at 950°C, and BaCO$_3$ from orthorhombic to hexagonal at 800°C (further to cubic at 970°C) [25]. On the other hand, the thermal decomposition temperature of the carbonates can be lowered drastically by adding certain oxides. For instance, the thermal decomposition temperature of SrCO$_3$ was lowered from 1150°C to 1102°C by adding a specific amount of Al$_2$O$_3$ [26]. At temperature higher than the SrCO$_3$ transformation temperature, 950°C, the solid-state reaction is dominated by diffusion of ions in the SrCO$_3$ lattice. Cations diffusion can be enhanced by the orthorhombic to hexagonal transformation of SrCO$_3$ because the high-temperature modification has a lower density. On the other hand a kinetic study of SrCO$_3$ decomposition in the solid-state reaction reveals that the activation energy of SrCO$_3$ decomposition was influenced by the calcination temperature rather than the nature of precursor. The interfacial reaction at low temperatures is characterized by a high activation energy of 130 kJ/mol, whereas, in the reactions at higher temperatures the activation energy of SrCO$_3$ decomposition decreases drastically to 34 kJ/mol [27]. MgCO$_3$ decomposed at 450-550°C and SrCO$_3$ above 1000°C so the dopants are released gradually in the system.
Thermal decomposition of hydromagnesite 4 \( \text{MgCO}_3 \cdot \text{Mg(OH)}_2 \cdot 5\text{H}_2\text{O} \) proceeds via dehydration (< 300°C) and decarbonation (>350°C) towards MgO. The decarbonation is highly dependent on the partial pressure of carbon dioxide [28] (fig1.a). From the figure 1.b it can be underlined that the decomposition of SrCO\(_3\) started at around 900°C and is finished at 1100°C. The endothermal event at 947°C is attributable to the phase transition of Sr carbonate (orthorhombic to hexagonal).

By controlling the conversion rate of strontium carbonate (decomposition) one could try favoring the incorporation of Sr\(^{2+}\) for La\(^{3+}\) sites in LSGM phase.

Considering the differential thermal analysis DTA/TG (thermogravimetry) results of stoichiometric \( \text{La}_{0.8}\text{Sr}_{0.2}\text{Ga}_{0.8}\text{Mg}_{0.2}\text{O}_{3-x} \) starting mixture with two heating rates, presented in figure 2, one can mention at least four weight loss distinct regions determined mainly by the elimination of adsorbed and crystalline water (up to 345°C), and decomposition of carbonates at 450-550°C for MgCO\(_3\) and 750-940°C for SrCO\(_3\) respectively. The lower temperature for the last reaction is explained by the influence of other components of the system (especially the presence of Mg\(^{2+}\) cations). Above 940°C no weight loss can be observed. It is worth to mention some endothermal events recorded on DTA curve at 849°C, 960°C, 1260°C, 1352°C, 1425°C and 1440°C respectively, assigned to the reactions which occurred in nonisothermal conditions when the heating rate was 5°C/min.

The phase evolution of the precursor with calcinations temperature at 1000°C for 6 h and samples sintered at 1450°C and 1500°C for 6 h is shown by means of XRD measurements (fig 3). It can be seen that the perovskite phase was identified in the pattern of the calcined samples, which is in agreement with the research results indicating that the perovskite phase transformation temperature is 1000°C. Other intermediate phases including \( \text{LaSrGa}_3\text{O}_7 \), \( \text{La}_4\text{Ga}_2\text{O}_9 \), and some trace of carbonates were also detected.

\( \text{LaSrGa}_3\text{O}_7 \) crystallizes in a tetragonal structure of melilite-type \( \text{ABC}_3\text{O}_7 \). \( \text{LaSrGa}_3\text{O}_7 \) belongs to space group \( \text{P}-421\text{m} \) with the lattice parameters \( a_0 = 8.056 \text{ Å} \) and \( c_0 = 5.333 \text{ Å} \) [29]. According to phase-diagram studies [17, 19, 20]
La$_{1-x}$Sr$_x$Ga$_2$O$_7$ is the most stable high-temperature phase Ga$_2$O$_3$–La$_2$O$_3$–SrO system. The phase has extended coexistence regions with the following compounds: Ga$_2$O$_3$, (La,Sr)GaO$_3$, SrGaO$_2$, Sr$_2$Ga$_2$O$_{13.5}$, LaSrGaO$_4$, La$_2$SrGa$_3$O$_{12}$, LaSr$_2$Ga$_5$O$_{15}$, and a liquid in the strontium rich part of the Ga$_2$O$_3$–SrO system [18]. The compound has a broad single-phase region at 1470°C in air (i.e., $x = 0.20$ to 0.60). When the lanthanum content is increased, the unit-cell volume decreases in size, particularly along the c-axis.

No trace of impurity phases, i.e. LaSrGa$_3$O$_7$, was detected in the sample sintered in the temperature range of 1450°C - 1500°C for 6 h, which involve that the LaSrGa$_3$O$_7$ phase was removed (transformed) in the synthesis process of the Sr- and Mg-doped LaGaO$_3$. The indexed phase belong to a primitive cubic unit cell with $a = 3.9049$ Å and $V_0 = 59.545$ Å$^3$, smaller than that reported by Huang et al. [21].

Further detailed researches have to be focused on this part of the subsolidus phase diagram of the doped La-Ga-O system.

It is well known that the powders synthesized through different synthesis routes have different properties, including morphology, particle size, specific surface area, etc., and so resulting in different sinterabilities. In fact, the loss of sinterability when increasing the calcinations temperature of the starting materials is well known among the ceramists. The powders used to prepare the green compacts were calcined at 1000°C for 6 h and so higher temperature reduced the starting powders sinterability. It can be observed that the shrinkage ratio at 1500°C was about 10%.

Figure 4.a presents the FTIR spectra of raw main materials. La$_2$O$_3$ spectrum presents absorption bands at 443, 643, 803, 1029, 1099, 1262, 1457 and 3609 cm$^{-1}$ respectively. The absorption bands 643 cm$^{-1}$ and 3609 cm$^{-1}$ correspond to hexagonal La(OH)$_3$, which forms as a consequence of moisture absorption from air. The peak at 1457 cm$^{-1}$, 1099 cm$^{-1}$, 1029 cm$^{-1}$ and 803 cm$^{-1}$ corresponds to the stretching vibration of the carbonate group. The absorption band of cubic phase La$_2$O$_3$ appears at 520 cm$^{-1}$ and 443 cm$^{-1}$. Prominent IR absorption bands are observed at 458 cm$^{-1}$, 511 cm$^{-1}$, 665 cm$^{-1}$, 698 cm$^{-1}$, 750 cm$^{-1}$ in the case of Ga$_2$O$_3$. The bands located at 458 cm$^{-1}$ and 511 cm$^{-1}$ could be assigned Ga–O stretching vibration modes. The shoulders at 750 cm$^{-1}$ was noticed by Burkhalder [30] in FTIR spectra of Ga$_2$O$_3$ but not assigned to a specific vibration. Fig. 4b presents the IR spectra of doping precursors. The strong absorption band of CO$_3^2$- ion was observed at about 1475 cm$^{-1}$ which is characteristic for SrCO$_3$. Strong
narrow absorption bands at about 855 cm\(^{-1}\) and 704 cm\(^{-1}\) for SrCO\(_3\), are assigned to be out of plane bending vibrations and in plane bending vibrations, respectively. Weak narrow absorption bands at about 1068 cm\(^{-1}\) due to the symmetric stretching vibrations were also detected.

The band located at 3448 cm\(^{-1}\) was assigned to OH\(^-\) group vibration in Mg precursor. The band at 1657 cm\(^{-1}\) was assigned to H\(_2\)O bending vibration of interlayer water in 4MgCO\(_3\)Mg(OH)\(_2\) 5H\(_2\)O. The \(v_3\) vibration of CO\(_3\)^{2-} in MgCO\(_3\) was split in two signals, i.e., 1483 cm\(^{-1}\) and 1421 cm\(^{-1}\), and the \(v_1\) vibration band of CO\(_3\)^{2-} was observed at 1120 cm\(^{-1}\). The absorption intensity below 550 cm\(^{-1}\) is probably attributable to octahedral magnesium.

Details about spectra of La\(_2\)O\(_3\)-Ga\(_2\)O\(_3\) system are limited to those reported by Zhang et al. [31] and Saine et al. [32]. They have reported the infrared spectrum of the orthorhombic LaGaO\(_3\) phase. The IR spectrum showed a total of 25 IR-active and 8 inactive modes for Pnma (no.62) form and 8 IR-active and 5 inactive modes for R3c form, respectively. [33] The undoped LaGaO\(_3\) undergoes a first-order phase transition from orthorhombic (Pbnm) to hexagonal–rhombohedral (R3c) at approximate temperature of 150°C. Slater et al. [11] have also shown that undoped LaGaO\(_3\) undergoes a low temperature phase transition and the La\(_{0.6}\)Sr\(_{0.1}\)Ga\(_{0.8}\)Mg\(_{0.2}\)O\(_{2.85}\) maintains monoclinic (I2/a) structure up to 1000°C.

The bands observed in the spectrum (fig.5) can be assigned, by comparison of spectral profile and band positions, to the modes noticed for other perovskites. Assignments are made on the basis of the relative band positions. FT-IR spectra are in agreement with the XRD observations. The vibrations up to 1000 cm\(^{-1}\) are characteristic to perovskite structures. After pre-sintering the presence of CO\(_2\) group’s vibrations and traces of adsorbed water were also noticed. After sintering no vibrations were identified in the wave number range of 1300-4000 cm\(^{-1}\).

SEM images present the surface morphology of the sintered ceramic (fig.6a). Well-crystallized grains of polyedral shape are formed in the micrometric range in the normal ceramic microstructure with a good density. A grain to grain connectivity can be seen with average particle size 2-5 \(\mu\)m as shown in the higher magnification micrographs (fig. 6b). The presence of open porosity can be also observed (fig. 6b). This observation is in good agreement with the electrical behavior of the ceramic. These microstructures also explain the measured values for densities, porosities and the large shrinkage (~10% after sintering at 1500°C/6h) of the sintered ceramics. Dilatometric tests on calcined samples confirmed that the active shrinkage processes, probably associated with cations interdiffusion, start at approximately 850-1000°C. The results show a dependence between heating rate and shrinkage. That is, higher shrinkage (densification) is achieved at lower heating rates. This results yield fruitful information for the optimization of the process for such ceramics.

Fig 7 shows the dilatometric curves of LSGM ceramic undergoing a thermal cycle consisting of continuous heating and cooling. The dilatometric curve shows linear thermal expansion characteristics in the investigated temperature range.
range. The calculated thermal expansion coefficient (TEC) is $9.13 \times 10^{-6}$ K$^{-1}$ in the range 200-800°C. The strain appearing (heating and cooling) indicates that the specimen length is change by the thermal cycle, which implies that the volume change has non-isotropic characteristics. It is well known that the elastic modulus and TEC have a major contribution to the stresses that can lead to cracking. If the difference in the TECs of surrounding components is large, its may crack during thermal cycling due to thermally induced stresses.

The sintered ceramics were characterized by impedance spectroscopy (Wayne Kerr 6440A) at frequencies in the range 20 Hz – 3MHz with $V_{ac} = 1$ V. The measurements were carried out in air from room temperature up to 500°C. In Fig. 8 the imaginary part of the impedance is plotted versus the real part of the impedance for the sample heated at 300°C and 500°C, respectively.

At 300°C the spectra reveal two semicircles, which can be assigned to the diffusion of O$^2-$ ion from vacancy to vacancy inside the grains (bulk process) and between two neighboring grains (grain boundary process). The semicircle from higher frequency corresponds to bulk process whereas the semicircle from lower frequency corresponds to grain boundary process.

To determine the dc-conductivity of the bulk and grain boundary, the semicircles of Nyquist plot in Fig. 8 were extrapolated to the intersections with the real axis, which gives the dc-resistances. In our case, the arcs are depressed with their centers displaced below the real axis; consequently the accuracy of simple extrapolation is not sufficient. The data have to be described with electrical equivalent circuits. It is not possible to fit the semicircle with the response of a simple RC - circuit (ideal semicircle with center on the real axis) because the bulk, as well as the grain boundary region exhibits a distribution of relaxation times (both centers below the real axis). This behavior can be described by a RC-circuit, in which the capacitor is exchanged by a "constant phase element" (CPE). The impedance of the CPE is given by the following equation:

$$ Z_{CPE} = A(i\omega)^{-b}, $$

where $\omega$ is the angular frequency and $A$ and $b \in [0,1]$ are frequency independent constants. An example of the result of a fit is shown in the Fig. 8 for sample measured at 300°C and 500°C. The fit (continuous line) describes not only the shape and absolute values of the measured data (open circles) but also the frequency dependence of the data.

From fitting data with the equivalent circuit depicted in figure 8 is observed that, at 300°C, the
grain boundary resistance ($R_{gb} \sim 6.4 \times 10^4 \, \Omega$) is considerable higher than the grain bulk component ($R_b \sim 3 \times 10^4 \, \Omega$).

At 500 °C the semicircle corresponding to the bulk contribution is out with respect to the frequency scale of instrumentation. The remained large semicircle corresponds to the grain boundary impedance $\sim 6 \times 10^2 \, \Omega$ and the low frequency spike inclined at $\sim 45^\circ$ is associated with diffusion of oxygen molecules at the contact.

Figure 9 shows the Arrhenius representation of total electrical conductivity of the sample sintered at 1500°C for 6 h. The lower values, compared with the published one [2,3] can be due to the microstructure of the sintered ceramic (grain boundary conduction effect).

4. Conclusion

A specific solid state preparation method was used to obtain Sr and Mg doped LaGaO$_3$ ceramics, starting from oxides and carbonates. Different heating rate for carbonates decomposition were used. The dopans were gradually released in the system.

Detailed structural investigation demonstrates the obtaining of a well-crystallized ceramics of cubic symmetry in the range of 1450°C-1500°C.

The structure and the morphology of the obtained ceramic may explain also the electrical behavior of the grain boundary resistance ($R_{gb}$) which is considerable higher than the grain bulk component ($R_b$).

Further investigation on the mechanism of reaction and the phase equilibrium of the doped LaGa-O system and also on sintering process are necessary.

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